



The lubrication mechanisms of textured surfaces

Tribology Today, Istanbul

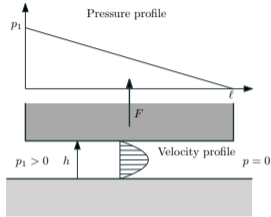
Noël Brunetière

April 18, 2026

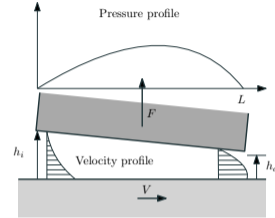
Institut Pprime, Poitiers, France

Fluid flow in a lubricated contact and load carrying capacity

Pressure flow (Hydrostatic effect)



Shear flow (Hydrodynamic effect)



- The mass flow rate per unit length is

$$\dot{m} = \frac{\rho}{12\mu} \frac{p_1}{\ell} h^3.$$

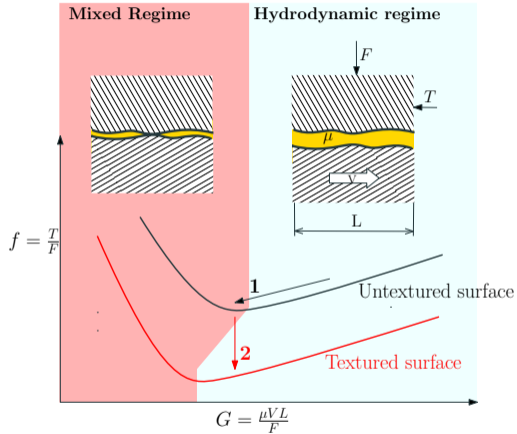
- Very strong impact of h on the flow.
- Constant force per unit length is $F = \frac{p_1}{2} \ell$.

- Force per unit length is

$$F = \frac{6\mu VL^2}{h_o^2} \left[\frac{\ln K}{(K-1)^2} - \frac{2}{K^2-1} \right] \text{ with } K = \frac{h_i}{h_o}$$

- Positive force can be generated depending on the slope. Zero force with $K = 1$ (flat on flat)

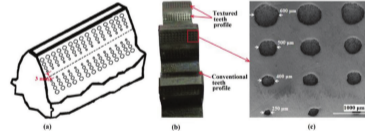
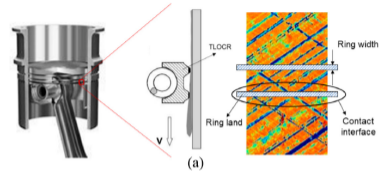
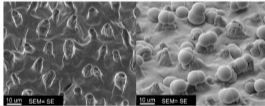
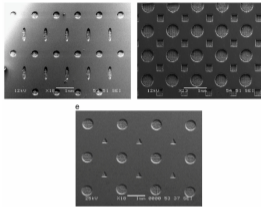
Lubrication regimes



Friction reduction

- Reduce G :
- Reduce size L or viscosity μ
- Risk of mixed lubrication and contact
- More particularly for flat on flat surfaces
- Surface texturing

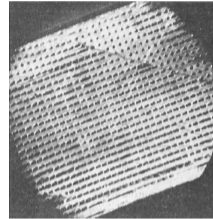
Surface texturing^{1,2}



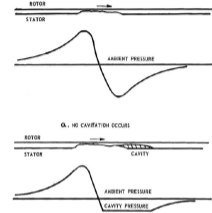
¹Philipp G. Grützmaier, Francisco J. Profito, and Andreas Rosenkranz. “**Multi-Scale Surface Texturing in Tribology-Current Knowledge and Future Perspectives**”. In: *Lubricants* 7.11 (2019). DOI: 10.3390/lubricants7110095.

²Andreas Rosenkranz et al. “**Surface Texturing in Machine Elements - A Critical Discussion for Rolling and Sliding Contacts**”. In: *Advanced Engineering Materials* 21.8 (2019), p. 1900194. DOI: 10.1002/adem.201900194.

- Origin of the load support in flat on flat lubrication
- Role of micro-irregularities (asperities)
- Experimental analysis of artificial rough surface
- Network of bumps by photo-etching
- Fluid film formation and load support due to surface texture



a)



b)

Figure 1: a) Surface texture used by Hamilton et al. and b) proposed load support mechanism

³D.B. Hamilton, J.A. Walowit, and C.M. Allen. “**A Theory of Lubrication by Micro-irregularities**”. In: *Journal of Basic Engineering* 88 (1966), pp. 177–185.

- No publication during about 30 years
- Development of surface texturing by laser and other technics in the 90'
- Increasing interest

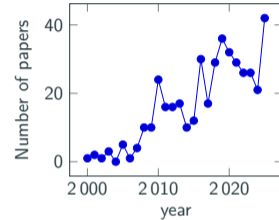
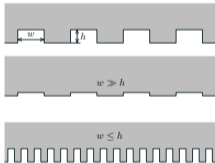


Figure 2: Number of papers: Google Scholar "allintitle: surface (texturing OR texture OR textured) (lubrication OR lubricated OR hydrodynamic)"

Example 1: shallow grooves⁴

- Friction of mechanical seals (flat on flat annular lubricated contact)
- texturing by ion-etching through a mask
- More than 50 % of friction reduction and no wear

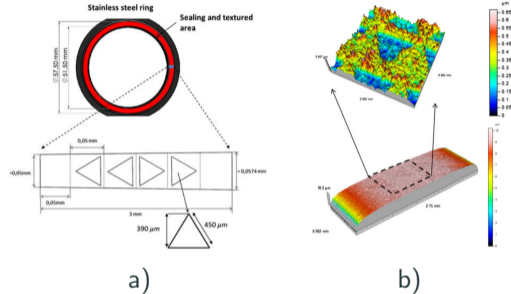
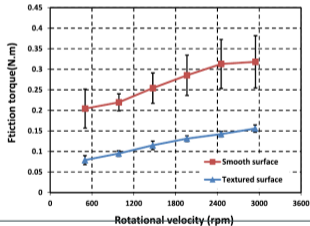


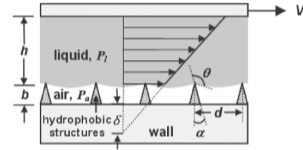
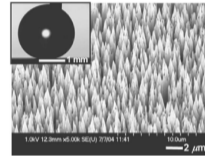
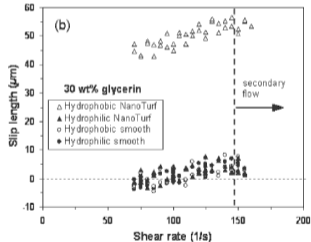
Figure 3: a) Configuration of the contact and surface texture and b) Topography of the surface

⁴M. Adjemout, N. Brunetière, and J. Bouyer. “Friction and temperature reduction in a mechanical face seal by a surface texturing: comparison between TEHD simulations and experiments”. In: *Tribology Transactions* 61.6 (2018), pp. 1084–1093. DOI:

10.1080/10402004.2018.1478053. URL: <https://doi.org/10.1080/10402004.2018.1478053>.

Example 2: deep grooves⁵

- Deep and narrow cavities / pillars
- Super-hydrophobic surface, poor wetting
- Allow apparent slip of the liquid, reduce friction



a)

b)

Figure 4: a) Surface texture and wetting b) Slip of the fluid

⁵Chang-Hwan Choi and Chang-Jin Kim. “Large Slip of Aqueous Liquid Flow over a Nanoengineered Superhydrophobic Surface”. In: *Phys. Rev. Lett.* 96 (6 2006), p. 066001. DOI: 10.1103/PhysRevLett.96.066001.

Outline of the course

1. Shallow grooves

Lubrication of one cavity

Multiple cavities in bearings

Multiple cavities in seals

Conclusions

2. Deep grooves

Wetting of surfaces

Effect of surface topography

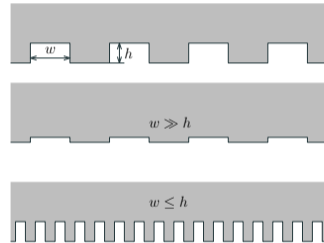
Flow over one cavity

Multiple cavities

Multiple cavities

Conclusions

3. Conclusions



Shallow grooves

Step bearing

- Lubrication assumptions ($h \ll L$)

- Reynolds equation:

$$\frac{\partial}{\partial x} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial x} \right) = \frac{U}{2} \frac{\partial \rho h}{\partial x}$$

$$\bar{h} = \frac{h_t}{h_f}$$

$$\text{Solution: } p_m = \frac{3\mu UL}{h^2} \times \frac{\bar{h}-1}{\bar{h}^3+1}$$

- Pressure $\propto L$

- h_t must be close to $2 \times h_f$

- No effect for deep grooves

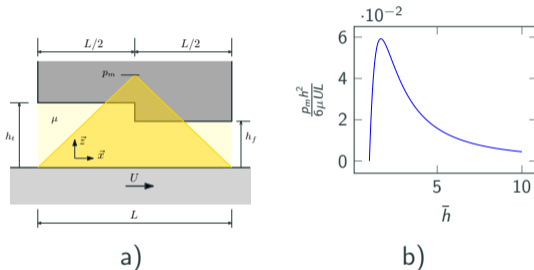
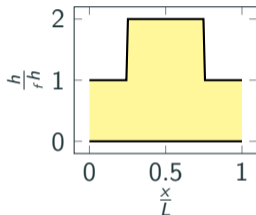


Figure 5: a) Configuration of a step bearing b) Max pressure

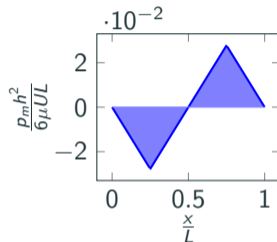
One centered cavity

- Lubrication assumptions ($h \ll L$)
- Reynolds equation:
 - $$\frac{\partial}{\partial x} \left(\frac{\rho h^3}{12\mu} \frac{\partial p}{\partial x} \right) = \frac{U}{2} \frac{\partial \rho h}{\partial x}$$

- Anti-symmetrical pressure distribution
- No load generation



a)



b)

Figure 6: a) Film thickness distribution b) Pressure distribution

One centered cavity with cavitation

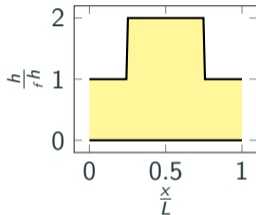
- Occurrence of cavitation:
- In cavitation zone: $p = p_c$
- In cavitation zone: $\rho \leq \rho_\ell$ where ρ_ℓ is the lubricant density

- No symmetry in the pressure distribution

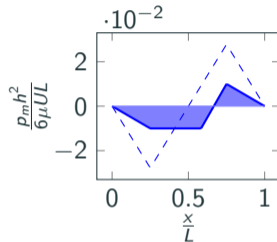
- $p_c = -0.01 \frac{6\mu UL}{h^2}$

- Negative load

- $p_{av} = \frac{1}{L} \int_0^L p dx = -0.0033 \times \frac{6\mu UL}{h^2}$



a)



b)

Figure 7: a) Film thickness distribution b) Pressure distribution

One cavity not centered

- The cavity is not in the center of the contact

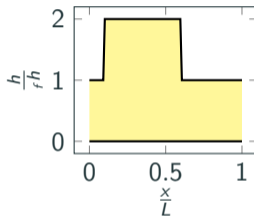
- Cavitation pressure

$$p_c = -0.01 \frac{6\mu UL}{h^2}$$

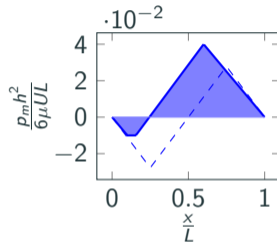
- No symmetry in the pressure distribution

- **Positive load**

- $p_{av} = -0.0136 \times \frac{6\mu UL}{h^2}$



a)



b)

Figure 8: a) Film thickness distribution b) Pressure distribution

Slider bearing⁶

- Slider bearing partially textured at the inlet
- Full texturing = one cavity, no load

- Pressure generation
- Less efficient than an equivalent step bearing
- Limited interest

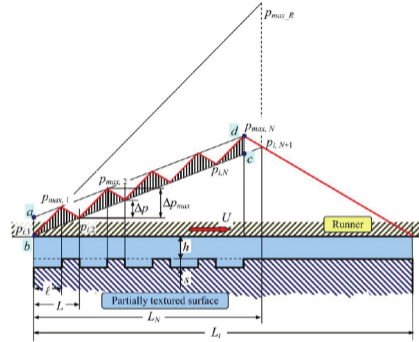


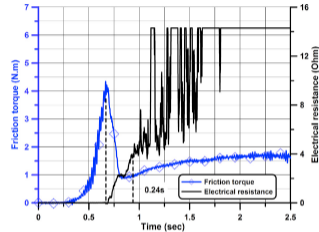
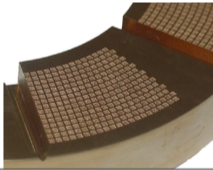
Figure 9: Configuration of the problem and pressure distribution

⁶M D Pascovici et al. “**Analytical investigation of a partially textured parallel slider**”. In: *Proceedings of the Institution of Mechanical Engineers, Part J: Journal of Engineering Tribology* 223.2 (2009), pp. 151–158. DOI: 10.1243/13506501JET470. URL: <https://doi.org/10.1243/13506501JET470>.

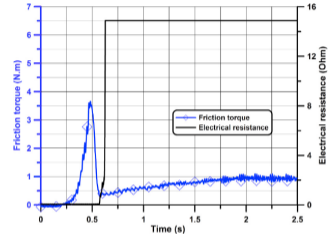
Thrust bearing⁷

- Thrust bearings with partial texturing and pocket

- Better efficiency of the pocketed bearing



a)



b)

Figure 10: Friction torque and contact resistance for a) textured bearing b) pocketed bearing

⁷Y. Henry, J. Bouyer, and M. Fillon. “**Experimental analysis of the hydrodynamic effect during start-up of fixed geometry thrust bearings**”. In: *Tribology International* 120 (2018), pp. 299–308. ISSN: 0301-679X. DOI: <https://doi.org/10.1016/j.triboint.2017.12.021>. URL: <https://www.sciencedirect.com/science/article/pii/S0301679X17305790>.

Principle of a mechanical seal

2 flat and smooth surfaces to close the gap

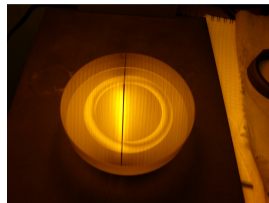
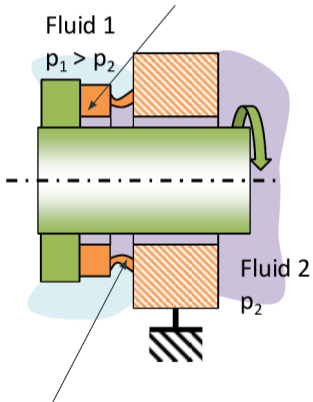


Figure 11: Surface flatness $\approx 0.3 \mu\text{m}$

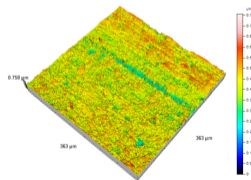


Figure 12: Surface roughness, $S_q = 0.05 \mu\text{m}$

Mechanical seal⁸

- Mechanical seal with one row of cavities
- Periodical boundary conditions

- Limited load generation
- Triangle provides negative load and significant radial flow

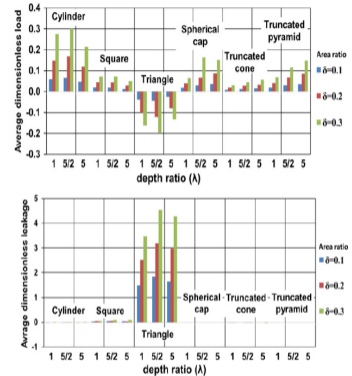
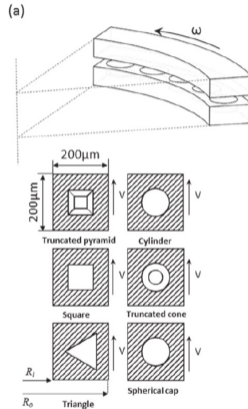


Figure 13: Load and leakage with the different configurations

⁸M Adjemout, N Brunetiere, and J Bouyer. “Numerical analysis of the texture effect on the hydrodynamic performance of a mechanical seal”. In: *Surface Topography: Metrology and Properties* 4.1 (2016), p. 014002. DOI: 10.1088/2051-672X/4/1/014002.

Example 1: shallow grooves¹⁰

- Friction of mechanical seals (flat on flat annular lubricated contact)
- texturing by ion-etching through a mask
- More than 50 % of friction reduction and no wear

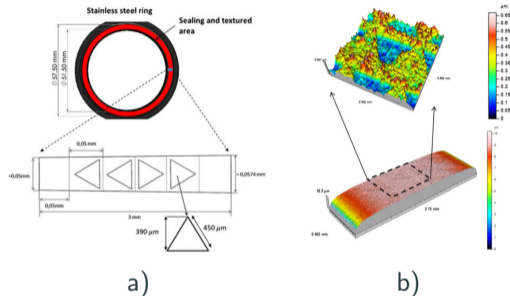
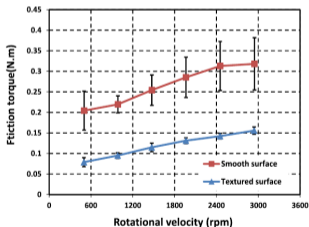


Figure 15: a) Configuration of the contact and surface texture and b) Topography of the surface

¹⁰Adjemout, Brunetière, and Bouyer, "Friction and temperature reduction in a mechanical face seal by a surface texturing: comparison between TEHD simulations and experiments".

Comparison to other texture types¹¹

- The other symmetrical texture should be bad !!
- Let's try !

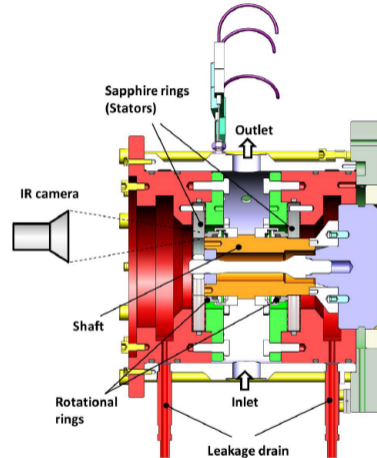
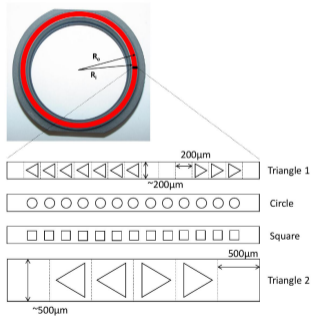


Figure 16: Test cell and temperature measurement

¹¹Kacou Dingui et al. “Surface Texturing to Reduce Temperature in Mechanical Seals”.

In: *Tribology Online* 15.4 (2020), pp. 222–229. DOI: 10.2474/trol.15.222.

Comparison to other texture types¹²

centering

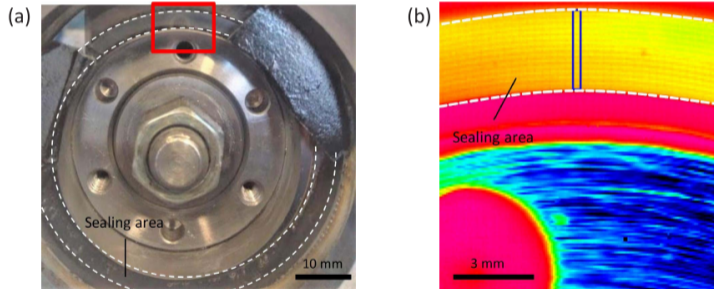


Fig. 5 a) View of the front of the test cell -b) Infrared view within the red square. The blue rectangle corresponds to the area analyzed to calculate the temperature.

¹²Dingui et al., "Surface Texturing to Reduce Temperature in Mechanical Seals".

Comparison to other texture types¹³

- Good results with other textures (symmetrical ones)
- Temperature controlled by texture area not the shape !!
- Why ??

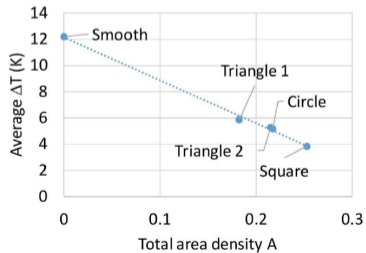
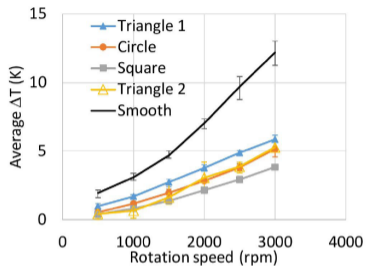
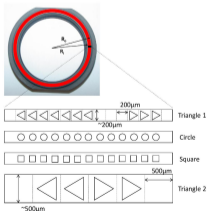


Figure 17: Main results

Why symmetrical texture works in seals?¹⁴

- Mechanical seal with 19 rows of spherical cap cavities
- Mixed lubrication model

- Rough surfaces can reach hydrodynamic regime
- Texturing improve this capability

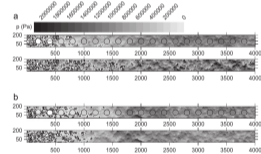
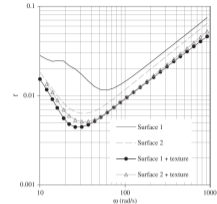
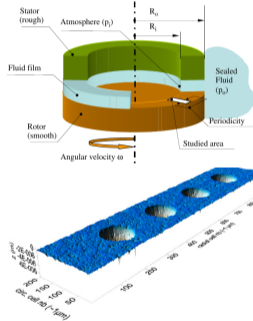


Figure 18: Stribeck curves and pressure distribution for rough and rough textured surfaces

¹⁴N. Brunetière and B. Tournerie. “Numerical Analysis of a Surface-Textured Mechanical Seal Operating in Mixed Lubrication Regime”. In: *Tribology International* 49 (2012), pp. 80–89. URL: <https://doi.org/10.1016/j.triboint.2012.01.003>.

Conclusion on shallow grooves

- Can surface texturing with shallow grooves help in flat on flat lubrication ?
- **Yes** by enhancing hydrodynamic pressure generation **but with conditions**
- Bearing like configurations:
 - Partial texturing located at the inlet only
 - However less efficient than a pocket (step bearing)
- Seal like configurations (flat ring on flat ring):
 - Triangular cavities placed base to base can provide a collective effect and load generation
 - No collective effect with symmetrical cavities
 - Symmetrical cavities can improve the natural pressure generation due to surface roughness
- Cavitation is not the mechanism that make surface texture working

Deep grooves

Contact angle: Energetic approach

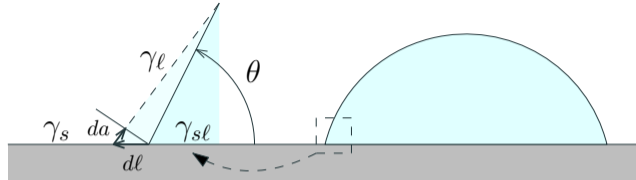


Figure 19: Energy balance at equilibrium

- Liquid surface energy variation $da\gamma_l$
- Solid surface energy variation $-dl\gamma_s$
- Solid liquid interfacial energy variation $dl\gamma_{sl}$
- At equilibrium the energy is minimized, small variation must vanish
- $da\gamma_l + dl\gamma_{sl} - dl\gamma_s = 0$
- Geometric relation: $da = \cos\theta dl$
- Young's equation: $\gamma_l \cos\theta = \gamma_s - \gamma_{sl}$

Different surface types

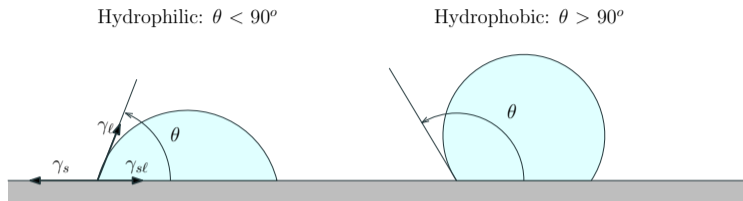


Figure 20: Surfaces classifications

- The angle characterizes the affinity of the solid and liquid;
- Small angle corresponds to a good affinity;
- High surface energy solids (high γ_s) are generally hydrophilic;
- Low surface energy solids (low γ_s) are generally hydrophobic;
- If $\gamma_s > \gamma_{sl} + \gamma_l$, the fluid totally spreads (no remaining drop).

Surface energy and wetting

- Examples of contact angle on solids of decreasing surface energy.

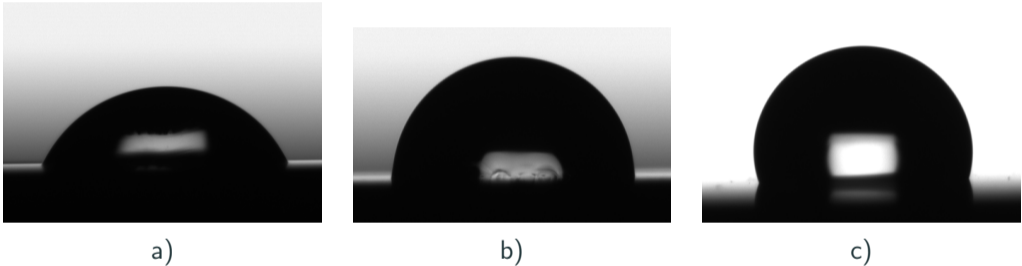


Figure 21: Sesille water drop on a) Stainless steel b) Silicon carbide c) PTFE

Wetting of a rough/textured surface

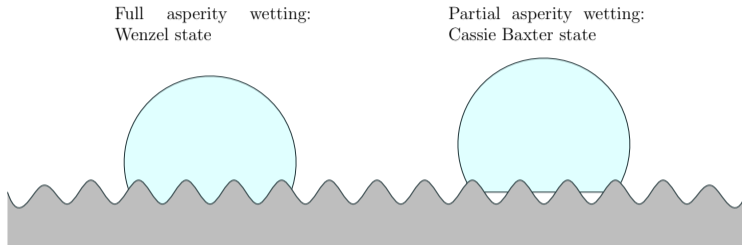


Figure 22: Different wetting regimes of a rough surfaces

- Different wetting regimes of a rough surface
- Wenzel state: the asperities are fully wetted
- Cassie-Baxter state (or fakir state): partial wetting of the asperities
- The most energetically favorable state will be reached

Wenzel regime

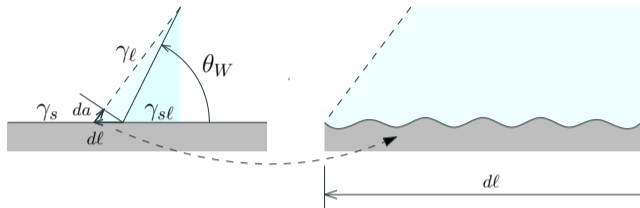


Figure 23: Energy balance at equilibrium

- The real length of the rough surface is linked to the linear length by $d_r l = R dl$
- R is a roughness factor ≥ 1 .
- Liquid surface energy variation $da \gamma_l$
- Solid surface energy variation $-d_r l \gamma_s$
- Solid liquid interfacial energy variation $d_r l \gamma_{sl}$
- Young equation: $\cos \theta_W \gamma_l = R (\gamma_s - \gamma_{sl}) = R \gamma_l \cos \theta$
- Wenzell relation: $\cos \theta_W = R \cos \theta$

Cassie-Baxter regime

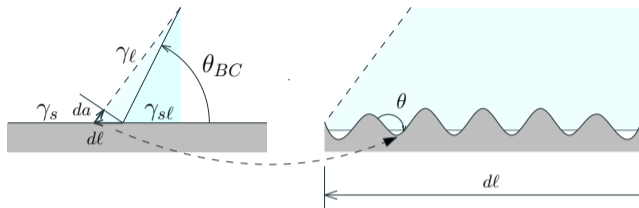


Figure 24: Energy balance at equilibrium

- f is the fraction of solid area wetted by the liquid (horizontally measured).
- Liquid surface energy variation $da\gamma_l$
- Solid surface energy variation $-d_r l\gamma_s$
- Solid liquid interfacial energy variation: $fd_r l\gamma_{sl} + (1-f)dl\gamma_l + (1-f)d_r l\gamma_s$
- Energy balance: $\cos\theta_{CB}\gamma_l + fR(\gamma_{sl} - \gamma_s) + (1-f)\gamma_l = 0$
- Cassie-Baxter relation: $\cos\theta_{CB} = fR\cos\theta + f - 1$

Example of contact angle on textured surfaces¹⁵

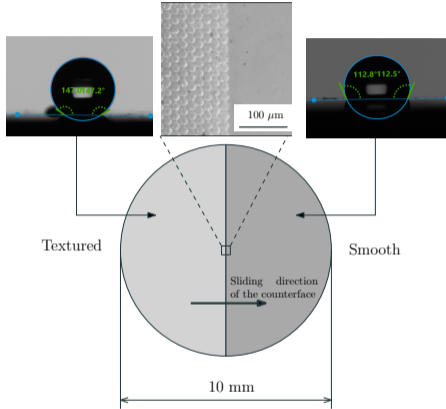


Figure 25: Increase of contact angle by surface texturing

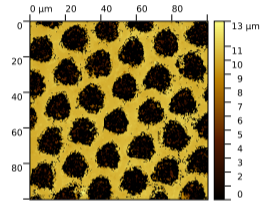
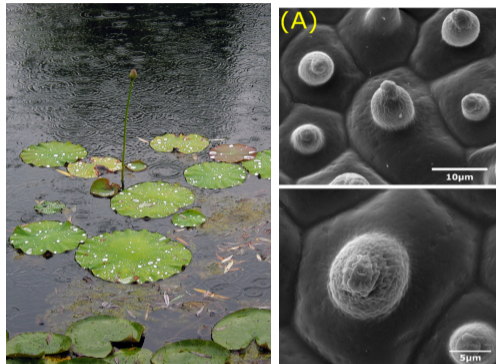


Figure 26: Topography of the rough surface

¹⁵Hiba Jendoubi, Olga Smerdova, and Noël Brunetière. **“Unexpected Frictional Behavior of Laser-Textured Hydrophobic Surfaces”**. In: *Lubricants* 9.3 (2021). ISSN: 2075-4442. DOI: 10.3390/lubricants9030031. URL: <https://www.mdpi.com/2075-4442/9/3/31>.

Natural superhydrophobic surfaces¹⁶



Lotus leaf

Low surface energy material combined with surface roughness leads to superhydrophobic behavior !

¹⁶Floris Honig et al. “**Natural Architectures for Tissue Engineering and Regenerative Medicine**”. en. In: *Journal of Functional Biomaterials* 11.3 (Sept. 2020), p. 47. DOI: 10.3390/jfb11030047. (Visited on 11/23/2020).

Wetting of a rough/textured surface

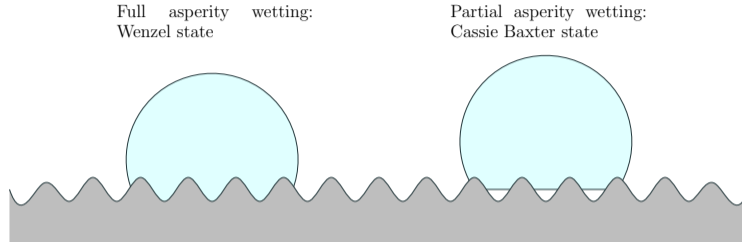
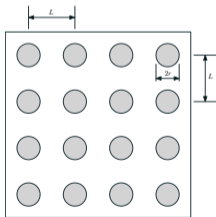


Figure 27: Different wetting regimes of a rough surfaces

- What will be the wetting regime ?
- The most energetically favorable state will be reached.
- It corresponds to the lowest contact angle of the two regimes:
- If $\theta_{CB} < \theta_W$, it will be the partial wetting;
- If $\theta_{CB} > \theta_W$, it will be the full wetting.

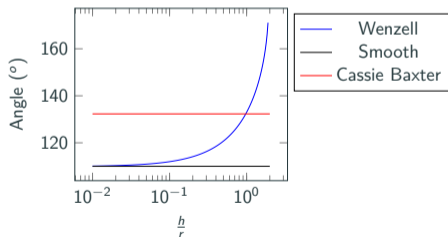
Wetting of a textured surface

- A flat plate with cylindrical cavities of depth h .
- Express the contact angle in the Cassie-Baxter and Wenzel regime
- Application to PTFE and water: $\theta = 110^\circ$, $r = 4\mu\text{m}$, $L = 10\mu\text{m}$



- $\cos \theta_{CB} = \left(1 - \frac{\pi r^2}{L^2}\right) \cos \theta - \frac{\pi r^2}{L^2}$
- $\cos \theta_W = \left(1 - \frac{2\pi r^2}{L^2} \times \frac{h}{r}\right) \cos \theta$

High contact angles require deep texture



Flow over one cavity^{17, 18}

- Navier Stokes solution (CFD - OpenFoam)
- Volume of Fluid + surface tension model
- Calculation of the flow, shear stress, pressure, etc
- Estimation of the apparent slip length
- Effect of the cavity size $\frac{d}{h_f}$ and $\frac{d}{h_c}$
- Effect of wetting (one phase flow or two phase flow)

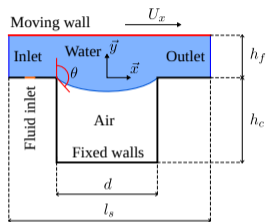


Figure 28: Configuration of the problem

¹⁷N. Elie. “Etude théorique et expérimentale des surfaces nano-texturées en régime de frottement lubrifié”. PhD thesis. Université de Poitiers, Dec. 2024.

¹⁸Nicolas Elie et al. “Slip Length in Shear Flow Over a Textured Surface”. In: *Tribology Online* 19.4 (2024), pp. 256–265. DOI: 10.2474/tro1.19.256. URL: <https://hal.science/hal-04664704>.

Apparent slip length

- Defined from the averaged speed profile
- Possibility to use the shear stress and the flow rate and compare to smooth channel:

- $Q = \frac{U_x h_f}{2}$ and $\tau = \frac{\mu U_x}{h_f}$

- Thus:

- $\frac{\delta_\tau}{h_f} = \frac{\mu U_x}{\tau_{av} h_f} - 1$

- $\frac{\delta_U}{h_f} = \frac{\frac{U_{av}}{U_x} - \frac{1}{2}}{1 - \frac{U_{av}}{U_x}}$

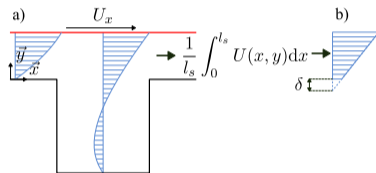


Figure 29: Usual slip length definition from average speed profile

Wenzel wetting regime

- Shear stress decreases when cavity depth increases h_c compared to d and when d increases compared to h_f
- 2 plateaus observed
- First one when $h_c \approx h_f$: formation of a full vortex in the cavity
- Second one when $h_c > d$: the location of the top vortex does not change
- Thus d must be $\leq h_c$ and much higher than film thickness h_f

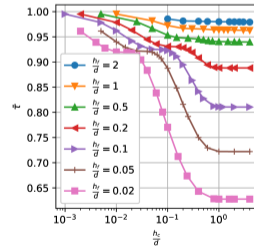
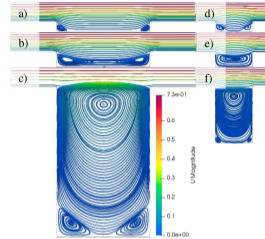
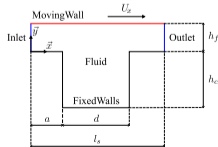


Figure 30: Stream lines for $\frac{hf}{d} = 0.2$ a-c) and $\frac{hf}{d} = 0.2$, Average shear stress

Wenzel wetting regime

- The slip length δ_τ based on the shear stress: monotonous evolution
- Slip length based on the average speed $\delta_U(U_{av})$ has a maximum value.
- Transition occurs when the top of the vortex changes from inside to outside the cavity.
- Thus hf/d must be maintained in the range 0.2–0.5.

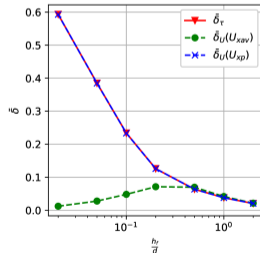
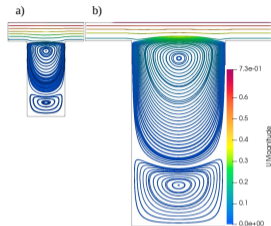
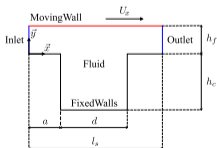


Figure 31: Stream lines for $\frac{hf}{d} = 0.2$ a-c) and $\frac{hf}{d} = 0.2$, slip length for $h_c = 2d$

Cassie Baxter wetting regime

- The viscosity difference leads to a shift of the vortex position / compared to full wetting
- Increasing d/h_f tends to center the vortex, as in mono-phase flow.

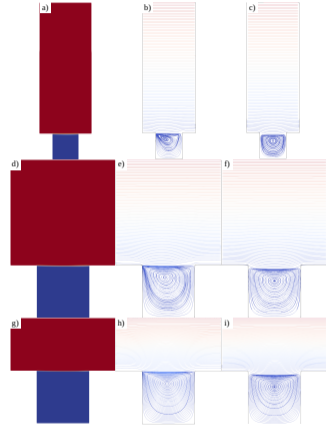
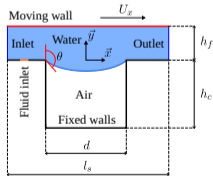


Figure 32: Stream lines in two-phase and mono-phase flow for $\frac{hf}{d} = 5$ a-c) $\frac{hf}{d} = 2$ d-f), $\frac{hf}{d} = 1$ h-j)

Cassie Baxter wetting regime

- $$\bar{\delta}_U = \frac{k_1 \frac{d}{h_f}}{k_2^2 + \frac{d}{h_f}^2}$$
- $$\bar{\delta} \approx \frac{k_1}{k_2^2} \frac{d}{h_f}, \text{ when } \frac{d}{h_f} \ll k_2 \text{ Consistent with literature}$$
- $$\bar{\delta} \approx k_1 \left(\frac{d}{h_f} \right)^{-1}, \text{ when } \frac{d}{h_f} \gg k_2.$$
- Higher slip with two-phase flow
- Better to have cavities with $\frac{d}{h_f} \approx k_2$

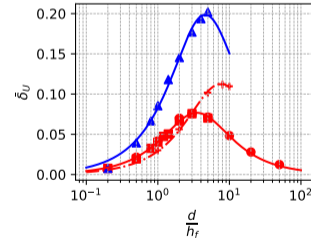
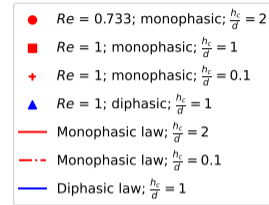
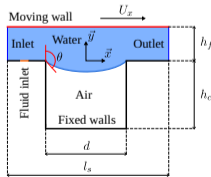


Figure 33: Slip length as a function of d/h_f for Wenzel and Cassie Baxter regimes

Effect of pressure on the liquid - gas interface

- Maximum pressure difference due to Laplace equation
- $p_{lmax} = p_g - \gamma \times \frac{2 \cos \theta}{d}$
- if $p_l > p_{lmax}$ the interface will enter the cavity
- Reducing cavity size d is better
- However $\delta \propto d$
- Example with water at 1 MPa in PTFE ($\theta = 110^\circ$), $d = 24 \text{ nm}$

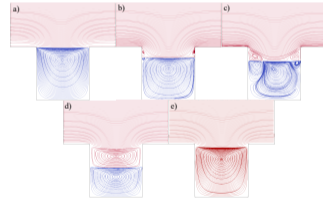
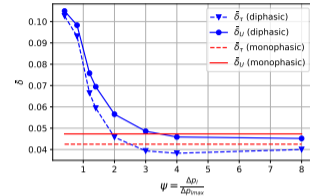
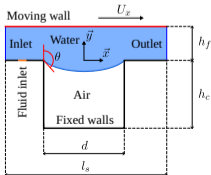


Figure 34: Streamlines for a size ratio $\frac{d}{h_f} = 1, 2$ and for the pressure ratios a) $\psi = 0.4$, b) $\psi = 3$, c) $\psi = 4$, d) $\psi = 8$ and e) monophasic. Evolution of slip length with pressure (top)

Slider bearing

- Slider bearing partially textured at the inlet
- Full texturing = no load

- Pressure generation
- Level of pressure sensitive to the wetting regime and cavity size

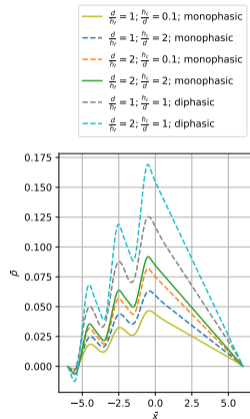
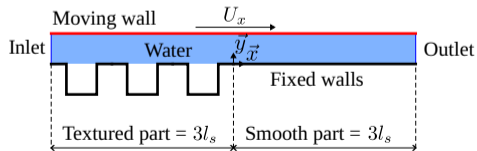


Figure 35: Pressure distribution for different wetting regimes and different cavity sizes

Slider bearing

- The pressure is controlled by the slip length
- Can be incorporated in the Reynolds equation:
- $\bar{p}_{max} = \frac{3\bar{\delta}_U}{2+5\bar{\delta}_U}$
- Allow to model deep cavities

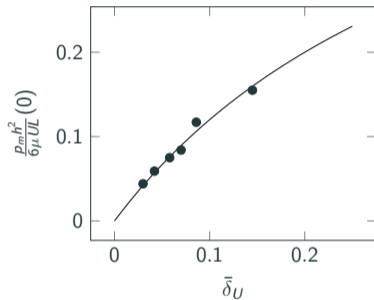
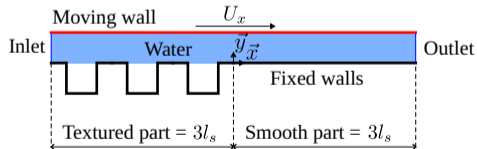


Figure 36: Pressure in the center of the slider as a function of the slip length

Conclusion on deep grooves

- Can surface texturing with deep grooves help in flat on flat lubrication ?
- **Yes** by creating an apparent slip of the liquid **but with conditions**
- Bearing like configurations:
 - Partial texturing located at the inlet only
 - Cavities must have a size with depth $h_c > d$, the width
 - The width d must be as wide as possible with a limit of about 2–4 times the film thickness h_f
 - Cassie Baxter is preferable but the liquid gas interface can enter the cavity at high pressure.

Conclusions

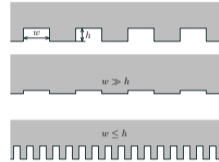
Conclusion on surface texturing in flat on flat lubrication

Shallow grooves $w \gg h$

- Mechanism is hydrodynamic lubrication
- Depth must be similar to the film thickness
- Inlet texturing for bearings
- Un-symmetrical texture for seals

Deep grooves $w < h$

- Mechanism is apparent slip
- Depth must be at least $1 \times$ width
- Width must be close to $2-4 \times$ film thickness
- Inlet texturing for bearings



- Both cases can be analyzed with Reynolds equation
- What happens between shallow and deep grooves?

Thank you for listening

Acknowledgments:

- Colleagues: Dr. J. Bouyer, Dr. P. Jolly, Pr. R. Lucas
- PhD students: Dr. M. Adjemout, Dr. N. Elie
- Companies: Latty, Cetim, IREIS
- Academic partners: Institut Jean Lamour, iRCER
- Projects:
 - Microgame ANR-2011-RMNP-00801
 - Labex Interatifs ANR-11-LABX-0017-01
 - EUR INTREE ANR-18-EURE-0010